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**Problem Sheet 2**  
**String Theory II**

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## Question 1. Torus-Partition Function and Modular Invariance: Free Boson

1. Fundamental Domain  $D$  of the torus: Let  $\tau \in \mathbb{C}$  be the modular parameter (modulus) of the torus  $T_\tau^2 = \mathbb{C}/\mathbb{Z} \oplus \tau\mathbb{Z}$ . Show that the torus  $T_\tau^2$  and  $T_{\gamma\tau}^2$ , where  $\gamma \in \mathrm{SL}_2(\mathbb{Z})$ , i.e.

$$\gamma\tau = \frac{a\tau + b}{c\tau + d}, \quad a, b, c, d \in \mathbb{Z}, \quad ad - bc = 1$$

describe the same space, i.e. the same identifications in  $\mathbb{C}$ . Using these  $\mathrm{SL}_2(\mathbb{Z})$  transformations one can restrict  $\tau$  to the fundamental domain

$$D : \quad -\frac{1}{2} \leq \operatorname{Re}(\tau) \leq \frac{1}{2}, \quad |\tau| \geq 1$$

Sketch  $D$  and determine the images of  $D$  under

$$\begin{aligned} T : \quad & \tau \rightarrow \tau + 1 \\ S : \quad & \tau \rightarrow -\frac{1}{\tau} \end{aligned}$$

2. For a single free boson  $X^\mu(z, \bar{z})$  in  $d$  dimensions, the torus partition function is

$$Z(\tau) = \operatorname{Tr} q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{c}{24}}$$

where  $q = e^{2\pi i\tau}$ . By performing the integral over momenta  $k$  and performing the sum over oscillator modes evaluate  $Z(\tau)$  and show that it is modular invariant, i.e.  $Z(\tau) = Z(\gamma\tau)$ , with  $\gamma \in \mathrm{SL}_2(\mathbb{Z})$ . You may use that  $\eta(\tau + 1) = e^{i\pi/12}\eta(\tau)$  and  $\eta(-1/\tau) = \sqrt{-i\tau}\eta(\tau)$ . What is the relevance of this result?

*Proof.* 1. Pick  $w_1, w_2 \in \mathbb{C}$  with  $\frac{w_1}{w_2} \notin \mathbb{R}$

The lattice  $\Lambda(w_1, w_2) := \{aw_1 + bw_2 : a, b \in \mathbb{Z}\}$

Tori  $\mathbb{C}/\Lambda(w_1, w_2) \cong \mathbb{C}/\Lambda(w'_1, w'_2) \iff \Lambda(w_1, w_2) = \Lambda(w'_1, w'_2)$

If  $\gamma \in \mathrm{SL}_2(\mathbb{Z})$ , then This exhausts the conformal symmetry?

$$\begin{pmatrix} w'_1 \\ w'_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} w_1 \\ w_2 \end{pmatrix} \Rightarrow w'_1, w'_2 \in \Lambda(w_1, w_2)$$

$$\begin{pmatrix} w'_1 \\ w'_2 \end{pmatrix} = \begin{pmatrix} d & -b \\ -c & a \end{pmatrix} \begin{pmatrix} w_1 \\ w_2 \end{pmatrix} \Rightarrow w_1, w_2 \in \Lambda(w'_1, w'_2)$$

$$\Rightarrow \Lambda(w_1, w_2) = \Lambda(w'_1, w'_2)$$

Conversely, if  $\gamma$  preserves the lattices, then

$$\det \gamma = ab - bc = 1 \Rightarrow \gamma \in \mathrm{SL}_2(\mathbb{Z})$$

$$\text{Let } \tau := \frac{w_1}{w_2}, \text{ then } \gamma(\tau) = \frac{a\tau + b}{c\tau + d}$$

This shows that the modular group is  $\mathrm{SL}_2(\mathbb{Z})$ .

(Should work in covariant quantisation  
& include the ghost sector)

□

2.  $Z(\tau) = \operatorname{Tr} (q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{c}{24}})$

where  $L_0 = \sum \alpha_n \cdot \alpha_0 + \sum_{n=1}^{\infty} \underbrace{\alpha_{-n} \cdot \alpha_n}_{\text{Number operator } N}$

$$\langle L_0 | m, \kappa \rangle = \alpha' \kappa^2 + m | m, \kappa \rangle$$

$$\operatorname{Tr} q^{L_0} = \prod_{n=1}^{\infty} \frac{1}{1 - q^n}$$

$$Z(\tau) = (q\bar{q})^{-\frac{c}{24}} \prod_{n=1}^{\infty} \frac{1}{1 - q^n} \prod_{m=1}^{\infty} \frac{1}{1 - \bar{q}^m} \int \frac{d\kappa}{2\pi} (q\bar{q})^{\kappa^2/2}$$

$$\propto \frac{1}{|\eta(\tau)|^2} \cdot \frac{1}{|\operatorname{Im} \tau|}$$

$$\begin{aligned} \eta\text{-function} : \quad & \eta(\tau+1) = e^{i\pi/12} \eta(\tau), \\ & \eta(-\frac{1}{\tau}) = \sqrt{-i\tau} \eta(\tau). \end{aligned}$$

Modular invariance ✓

## Question 2. Torus-Partition Function: RNS String

The one-loop or torus partition function for the closed RNS string is

$$Z_{T^2} = V_{10} \int_D \frac{d^2\tau}{2\tau_2} \int \frac{d^{10}k}{(2\pi)^{10}} \text{Tr}_{\mathcal{H}_k} (-1)^F q^{\alpha'(k^2 + M^2)/4} \bar{q}^{\alpha'(\bar{k}^2 + \bar{M}^2)/4},$$

where  $\mathcal{H}_k$  is the physical state (including GSO-projection) space with momentum  $k$  ground state, and  $q = e^{2\pi i \tau}$ , where  $\tau$  is again the modular parameter (modulus) of the torus  $T_\tau^2 = \mathbb{C}/\mathbb{Z} \oplus \tau\mathbb{Z}$  and  $D$  the fundamental domain. The spacetime Fermion number operator is  $F$  (not to be confused with the worldsheet one  $F$ ).

1. Evaluate the NS-sector and R-sector partition functions.
2. Using the results on theta-functions from the lecture, show that this partition function vanishes.

*Proof.* 1. (GSO-projection : See BLT §9.1)

$$Z_{RNS} = V_{10} \int_D \frac{d^2\tau}{2\tau_2} \int \frac{d^{10}k}{(2\pi)^{10}} \text{Tr}_{\mathcal{H}_k} (-1)^F q^{\alpha'(k^2 + M^2)/4} \bar{q}^{\alpha'(\bar{k}^2 + \bar{M}^2)/4}$$

GSO projector :  $P_{NS}^{GSO} = \frac{1}{2}(1 + (-1)^F)$   
 $P_R^{GSO} = \frac{1}{2}(1 + \eta(-1)^F)$        $F$ : chirality

States :

$$NS \quad \alpha_{-n_1}^{i_1} \cdots \alpha_{-n_N}^{i_N} b_{-r_1}^{j_1} \cdots b_{-r_M}^{j_M} |0\rangle_{NS}$$

$$R \quad \alpha_{-n_1}^{i_1} \cdots \alpha_{-n_N}^{i_N} b_{-r_1}^{j_1} \cdots b_{-r_M}^{j_M} |0\rangle_R \text{ (or } |\bar{0}\rangle_R)$$

Mass :  $M^2 = \frac{2}{\alpha'}(N_{tr}^{(a)} + N_{tr}^{(b)} + a) = \frac{2}{\alpha'}(\tilde{N}_{tr}^{(a)} + \tilde{N}_{tr}^{(b)} + a)$   
 $\hookrightarrow$  transverse

$$a_{NS} = -\frac{1}{2}, \quad a_R = 0$$

R-sector :

$$\frac{1}{2}(-1)^{(g+g)} \sum_n \sum_m \sum_r \sum_s \langle m, n, r, s | q^{N_{tr}^{(a)}} \bar{q}^{\tilde{N}_{tr}^{(a)}} q^{N_{tr}^{(b)}} \bar{q}^{\tilde{N}_{tr}^{(b)}} | m, n, r, s \rangle$$

$\underbrace{(-1)^F}_{\text{chirality}}$      $\underbrace{\alpha}_{\tilde{\alpha}}$      $\underbrace{b}_{\tilde{b}}$

degeneracy

$$\bullet \text{Bosons} : \sum_{m=0}^{\infty} q^{\sum_n \alpha_{-n} \cdot \alpha_n} = \sum_m d(m) q^m = \prod_{n=0}^{\infty} \frac{1}{1-q^n}$$

$$\bullet \text{Fermions} : \prod_{m=0}^1 \prod_r q^{b_{-r} \cdot b_r} = \prod_r (1-q^r)$$

$$\Rightarrow -8 \prod_i \left( \frac{1+q^i}{1-q^i} \right)^2 \prod_j \left( \frac{1+\bar{q}^j}{1-\bar{q}^j} \right)^2 \quad (\text{only contribution from } R)$$

NS-sector : No degeneracy in vacuum

$$\frac{1}{2} \prod_i \left( \frac{1+q^i}{1-q^i} \right)^2 \prod_j \left( \frac{1+\bar{q}^j}{1-\bar{q}^j} \right)^2 \cdot \frac{1}{q^{1/2}} \frac{1}{\bar{q}^{1/2}} - \frac{1}{2} \prod_i \left( \frac{1-q^i}{1-q^i} \right)^2 \prod_j \left( \frac{1-\bar{q}^j}{1-\bar{q}^j} \right)^2 \cdot \frac{1}{q^{1/2}} \frac{1}{\bar{q}^{1/2}}$$

$\frac{1}{2} (1 + \dots)$                                    $\frac{1}{2} (\dots + \underline{(-1)^F})$

### Question 3. Ghosts!

Let  $b$  and  $c$  be anticommuting fields,  $\beta$  and  $\gamma$  commuting fields, with action

$$S = \frac{1}{2\pi} \int d^2z (b\bar{\partial}c + \beta\bar{\partial}\gamma).$$

The OPE algebra is

$$b(z)c(0) \sim \frac{1}{z}, \quad c(z)b(0) \sim \frac{1}{z}, \quad \beta(z)\gamma(0) \sim -\frac{1}{z}, \quad \gamma(z)\beta(0) \sim \frac{1}{z}.$$

Define

$$\begin{aligned} T_{\text{ghost}}(z) &= (\partial b)c - \lambda\partial(bc) + (\partial\beta)\gamma - \frac{1}{2}(2\lambda - 1)\partial(\beta\gamma) \\ J_{\text{ghost}}(z) &= -\frac{1}{2}(\partial\beta)c + \frac{2\lambda - 1}{2}\partial(\beta c) - 2b\gamma \end{aligned}$$

1. Compute the conformal weights of  $b$  and  $c, \beta$  and  $\gamma$ .
2. Furthermore compute the OPE of  $TT$  and determine the central charge.
3. This system of ghosts can be used in the Faddeev-Popov gauge-fixing for  $\lambda = 2$ . Check that for this value the total central charge of  $T_{\text{ghost}}$  and  $T_{RNS}$  vanishes for  $d = 10$ .

*Proof.* 1.

□