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Problem Sheet 2
C3.11: Riemannian Geometry

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Section A: Introductory

Question 1

Let X, Y be vector fields on (M, g) . Let $p \in M$ and let $\alpha : (-\epsilon, \epsilon) \rightarrow M$ be the integral curve of X with $\alpha(0) = p$. For all $t \in (-\epsilon, \epsilon)$ let $\tau_t : T_p M \rightarrow T_{\alpha(t)} M$ be parallel transport along $\alpha|_{[0,t]}$. Show that

$$\nabla_X Y(p) = \left. \frac{d}{dt} (\tau_t^{-1} (Y(\alpha(t)))) \right|_{t=0}$$

Proof. We adopt Einstein's convention. A dot always denotes the derivative with respect to t , the parameter of a curve.

Choose a coordinate chart $(U; x^1, \dots, x^n)$ at p such that $x^i(p) = 0$. Without loss of generality suppose that $\alpha(-\epsilon, \epsilon) \subseteq U$. $\{\partial_1, \dots, \partial_n\}$ is a basis of $T_p M$. We parallelly transport the basis vectors along α . Let $E_i(t) := \tau_t(\partial_i)$. Then $\{E_1, \dots, E_n\}$ is a basis of parallel vector fields along α . Suppose that $Y(\alpha(t)) = Y^i(t)E_i(t)$. Since α is the integral curve of X , $X = \dot{\alpha}$ on α . Then

$$\nabla_X Y(p) = \nabla_{\dot{\alpha}} Y(\alpha(t))|_{t=0} = \dot{\alpha}(Y^i(t))E_i(t)|_{t=0} = \dot{Y}^i(0)\partial_i$$

On the other hand,

$$\left. \frac{d}{dt} (\tau_t^{-1} (Y(\alpha(t)))) \right|_{t=0} = \left. \frac{d}{dt} (\tau_t^{-1} (Y^i(t)E_i(t))) \right|_{t=0} = \left. \frac{d}{dt} (Y^i(t)\partial_i) \right|_{t=0} = \dot{Y}^i(0)\partial_i$$

The result follows from above. □

Question 2

Let (M, g) be an n -dimensional Riemannian manifold. Let $p \in M$ and let U be a normal neighbourhood of p . Let $\{E_1, \dots, E_n\}$ be an orthonormal basis for $T_p M$, let $\psi : T_p M \rightarrow \mathbb{R}^n$ be given by $\psi(\sum_{i=1}^n x_i E_i) = (x_1, \dots, x_n)$ and let $\varphi = \psi \circ \exp_p^{-1} : U \rightarrow \mathbb{R}^n$.

(a) Let $\gamma(t)$ be a geodesic through p in U in M . Show that

$$\varphi \circ \gamma(t) = (a_1 t, \dots, a_n t)$$

for $(a_1, \dots, a_n) \in \mathbb{R}^n$.

(b) Show that in (U, φ) , we have $g_{ij}(p) = \delta_{ij}$ and $\Gamma_{ij}^k(p) = 0$.

(c) Hence, or otherwise, show that there is open set $V \ni p$ and orthonormal vector fields E_1, \dots, E_n on V such that

$$\nabla_{E_i} E_j(p) = 0$$

Proof. (a)

□

Section B: Core

Question 3

Let (M, g) be a Riemannian manifold. Recall that a **Killing field** on M is a vector field X such that $\mathcal{L}_X g = 0$ or, equivalently, that the flow of X near any point consists of local isometries.

- (a) Let $p \in M$ and let U be a normal neighbourhood of p . Suppose that X is a Killing field on (M, g) so that $X(p) = 0$ and $X(q) \neq 0$ for all $q \in U \setminus \{p\}$.

By using the First variation formula, or otherwise, show that X is tangent to all sufficiently small geodesic spheres centred at p .

- (b) Show that X is a Killing field on (M, g) if and only if, for all vector fields Y, Z on M ,

$$g(\nabla_Y X, Z) + g(\nabla_Z X, Y) = 0$$

Proof. A dot always denotes the derivative with respect to t , the parameter of a curve.

- (a) Let $\gamma : [0, L] \rightarrow U$ be a radial geodesic starting from p . By Proposition 3.15, there exists a variation f of γ such that $X = X_f$ satisfies $X_f(t) = \partial_s f(0, t)$ along γ . The first variation formula is given by

$$\frac{1}{2}E'_f(0) = - \int_0^L g(X_f, \nabla_{\dot{\gamma}} \dot{\gamma}) dt + g(X_f(t), \dot{\gamma}(t)) \Big|_{t=0}^{t=L},$$

where the energy¹ is given by

$$E_f(s) = \int_0^L g(\dot{f}(s, t), \dot{f}(s, t)) dt.$$

Note that

$$\frac{d}{ds} g(\dot{f}(s, t), \dot{f}(s, t)) \Big|_{s=0} = (\mathcal{L}_X g)(\dot{\gamma}(t), \dot{\gamma}(t)) = 0,$$

since X is a Killing vector field. Hence

$$E'_f(0) = \int_0^L (\mathcal{L}_X g)(\dot{\gamma}(t), \dot{\gamma}(t)) dt = 0.$$

Since γ is a geodesic, $\nabla_{\dot{\gamma}} \dot{\gamma} = 0$. We also have $X_f(0) = X(p) = 0$ by assumption. Substituting into the first variation formula, we have

$$g(X_f(L), \dot{\gamma}(L)) = 0.$$

X is orthogonal to γ at $t = L$. Since γ is orthogonal to the geodesic spheres centred at p , which have tangent spaces of codimension 1, then X is tangent to the the geodesic spheres centred at p . ✓

- (b) For any $Y, Z \in \Gamma(TM)$, $g(Y, Z) = \text{tr}(g \otimes Y \otimes Z)$, where tr denotes the contraction of *all* covariant and contravariant indices of $g \otimes Y \otimes Z \in \Gamma(T_2^2 M)$. Since tr commutes with the Lie derivatives, we have

$$\begin{aligned} \mathcal{L}_X(g(Y, Z)) &= \mathcal{L}_X(\text{tr}(g \otimes Y \otimes Z)) = \text{tr}(\mathcal{L}_X(g \otimes Y \otimes Z)) \\ &= \text{tr}((\mathcal{L}_X g) \otimes Y \otimes Z) + \text{tr}(g \otimes \mathcal{L}_X Y \otimes Z) + \text{tr}(g \otimes Y \otimes \mathcal{L}_X Z) \\ &= (\mathcal{L}_X g)(Y, Z) + g([X, Y], Z) + g(Y, [X, Z]). \end{aligned}$$

¹**Cultural Remark.** For a free particle moving along the worldline γ on a Riemannian/Lorentzian manifold M , the Lagrangian is given by $\mathcal{L}[(\gamma(t))] = \frac{1}{2}mg(\dot{\gamma}, \dot{\gamma})$. The energy defined in the lecture should really be called the **action** $S[\gamma] = \int_0^{t_0} \mathcal{L}[\gamma(t)] dt$.

On the other hand, we have

$$\begin{aligned}\mathcal{L}_X(g(Y, Z)) &= X(g(Y, Z)) \\ &= g(\nabla_X Y, Z) + g(Y, \nabla_X Z) \\ &= g(\nabla_Y X, Z) + g(Y, \nabla_Z X) + g([X, Y], Z) + g(Y, [X, Z]).\end{aligned}$$

Comparing the two equations, we obtain

$$(\mathcal{L}_X g)(Y, Z) = g(\nabla_Y X, Z) + g(Y, \nabla_Z X)$$

In particular, $\mathcal{L}_X g = 0$ if and only if $g(\nabla_Y X, Z) + g(Y, \nabla_Z X) = 0$ for all $Y, Z \in \Gamma(TM)$. □

Perfect.

Question 4

Let (M, g) be a Riemannian manifold, let $f : M \rightarrow \mathbb{R}$ be a smooth function and let X be a vector field on M .

- (a) Note that we have a linear map from vector fields to vector fields given by $Y \mapsto \nabla_Y X$. We define the **divergence** of X to be the smooth function

$$\operatorname{div} X = \operatorname{tr}(Y \mapsto \nabla_Y X)$$

Show that if X is a Killing field then $\operatorname{div} X = 0$.

- (b) Recall that $Y \mapsto g(Y, -)$ defines an isomorphism between vector fields and 1-forms on M . We define the **gradient** of f to be the vector field ∇f given by

$$g(\nabla f, -) = df.$$

We define the **Laplacian** of f to be the smooth function

$$\Delta f = \operatorname{div} \nabla f$$

Show that

$$\Delta(f^2) = 2f\Delta f + 2|\nabla f|^2.$$

Now suppose further that M is compact, connected and oriented with Riemannian volume form Ω .

- (c) Show that

$$\mathcal{L}_X \Omega = (\operatorname{div} X)\Omega$$

Relate this to the result about Killing fields from (a).

- (d) Show that if $\Delta f \geq 0$ on M then f is constant.

Proof. We adopt Einstein's convention. Let us work in coordinates! Let $(U; x^1, \dots, x^n)$ be a chart and $\{\partial_1, \dots, \partial_n\}$ be the frame vector fields.

only if the frame is orthonormal at the point p you're considering, which of course you can find

- (a) Then $\operatorname{div} X = \nabla_\mu X^\mu \bigcirc g(\partial_\mu, \nabla_\mu X)$. Since X is a Killing vector field, from Question 3.(b), we have

$$(\mathcal{L}_X g)(\partial_\mu, \partial_\mu) = 2g(\partial_\mu, \nabla_\mu X) = 2 \operatorname{div} X = 0.$$

Hence $\operatorname{div} X = 0$. ✓

- (b) The musical isomorphism $TM \cong T^*M$ identifies the frame vectors ∂_μ with frame covectors dx^μ . So $\nabla f = \frac{\partial f}{\partial x^\mu} \partial_\mu$. Then

$$\nabla(f^2) = \frac{\partial(f^2)}{\partial x^\mu} \partial_\mu = 2f \frac{\partial f}{\partial x^\mu} \partial_\mu = 2f \nabla f$$

And

$$\Delta(f^2) = \operatorname{div}(2f \nabla f) = 2 \nabla_\mu (f \nabla f)^\mu = 2 \partial_\mu f (\nabla f)^\mu + 2f (\nabla_\mu \nabla f)^\mu = 2g(\nabla f, \nabla f) + 2f \Delta f \quad \checkmark$$

- (c) Let $p \in M$. We take $(U; x^1, \dots, x^n)$ to be the geodesic normal coordinates at p , such that

$$x^\mu(p) = 0, \quad g_{\mu\nu}(p) = \delta_{\mu\nu}, \quad \Gamma_{\mu\nu}^\lambda(p) = 0.$$

The (pull-back of) volume form at p in the coordinates is given by

$$\Omega = \sqrt{\det(\varphi^*g)} dx^1 \wedge \dots \wedge dx^n = dx^1 \wedge \dots \wedge dx^n.$$

By Cartan's formula,

$$\mathcal{L}_X \Omega = d \circ \iota_X \Omega + \iota_X \circ d\Omega = d \circ \iota_X \Omega.$$

If $X = X^\mu \partial_\mu$ at p , then we can compute

$$\begin{aligned} \iota_X \Omega &= \iota_X (dx^1 \wedge \dots \wedge dx^n) \\ &= (-1)^{\mu+1} \iota_X (dx^\mu) dx^1 \wedge \dots \wedge \widehat{dx^\mu} \wedge \dots \wedge dx^n \\ &= (-1)^{\mu+1} X^\mu dx^1 \wedge \dots \wedge \widehat{dx^\mu} \wedge \dots \wedge dx^n. \end{aligned}$$

Then

$$\begin{aligned} d \circ \iota_X \Omega &= (-1)^{\mu+1} dX^\mu \wedge dx^1 \wedge \dots \wedge \widehat{dx^\mu} \wedge \dots \wedge dx^n + (-1)^{\mu+1} X^\mu d(dx^1 \wedge \dots \wedge \widehat{dx^\mu} \wedge \dots \wedge dx^n) \\ &= (-1)^{\mu+1} \partial_\nu X^\mu dx^\nu \wedge dx^1 \wedge \dots \wedge \widehat{dx^\mu} \wedge \dots \wedge dx^n \\ &= \partial_\mu X^\mu dx^1 \wedge \dots \wedge dx^n \\ &= (\operatorname{div} X) \Omega. \end{aligned}$$

No Christoffel symbols appeared in the calculation because everything was evaluated at p . Therefore we obtain $\mathcal{L}_X \Omega = (\operatorname{div} X) \Omega$ at p . Since p is arbitrary, the relation holds on all M . \checkmark

- (d) We need to assume that M has no boundary.

By (c), we can consider the integral

$$\int_M \Delta f \Omega = \int_M \operatorname{div} \nabla f \Omega = \int_M \mathcal{L}_{\nabla f} \Omega = \int_M d \circ i_{\nabla f} \Omega = \oint_{\partial M} i_{\nabla f} \Omega = 0.$$

Since $\Delta f \geq 0$ on M , we must have $\Delta f = 0$ on M . By (b) we have

$$\Delta(f^2) = 2g(\nabla f, \nabla f) \geq 0,$$

because g is positive definite. Then following the same argument we must have $\Delta(f^2) = 0$ on M . Hence $g(\nabla f, \nabla f) = 0$ on M . By definiteness of g , we have $\nabla f = 0$. Since M is connected, f is constant on M . \checkmark *Perfect!* \square

Section C: Optional

Question 5

The Euclidean Schwarzschild metric (of mass $m > 0$) is defined for $(\cos \frac{t}{4m}, \sin \frac{t}{4m}) \in \mathcal{S}^1$, $r > 2m$, $\theta \in (0, \pi)$ and $(\cos \phi, \sin \phi) \in \mathcal{S}^1$ by

$$g = \left(1 - \frac{2m}{r}\right) dt^2 + \left(1 - \frac{2m}{r}\right)^{-1} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

and extends smoothly to $\theta = 0, \pi$.

- (a) Show that there are no geodesics in this metric with r constant.
- (b) Show that, given any point p with $r > 2m$ there exists a finite length geodesic γ starting at p ending at a point q with $r = 2m$.

Proof. (a) We can derive the radial equation for a geodesic exactly the same way as in General Relativity (where the Schwarzschild metric is Lorentzian).

The Lagrangian is given by

$$\mathcal{L}[\gamma(s)] = \left(1 - \frac{2m}{r}\right) \dot{t}^2 + \left(1 - \frac{2m}{r}\right)^{-1} \dot{r}^2 + r^2 (\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2),$$

where the dot denotes the derivative with respect to the affine parameter s . We observe that t and φ are ignorable coordinates. We have

$$\frac{\partial \mathcal{L}}{\partial \dot{t}} = 2 \left(1 - \frac{2m}{r}\right) \dot{t} = \text{const}, \quad \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = 2r^2 \sin^2 \theta \dot{\phi} = \text{const}.$$

We can use the $\text{SO}(3)$ symmetry of the manifold to fix the geodesics on the plane $\theta = \pi/2$. Then $\dot{\theta} = 0$. We set the constants $J := r^2 \dot{\phi}$ and $E := \left(1 - \frac{2m}{r}\right) \dot{t}$, which are the angular momentum and energy per unit mass.

Since γ is affinely parametrised, we have $\mathcal{L}[\gamma(s)] = g(\dot{\gamma}, \dot{\gamma}) = 1$. This gives

$$\mathcal{L} = \left(1 - \frac{2m}{r}\right) \dot{t}^2 + \left(1 - \frac{2m}{r}\right)^{-1} \dot{r}^2 + r^2 \dot{\phi}^2 = \left(1 - \frac{2m}{r}\right)^{-1} E^2 + \left(1 - \frac{2m}{r}\right)^{-1} \dot{r}^2 + \frac{J^2}{r^2} = 1.$$

Suppose that there is a geodesic with $r = \text{const}$. Then

$$E^2 = \left(1 - \frac{J^2}{r^2}\right) \left(1 - \frac{2m}{r}\right)$$

□